

#### Available online at www.sciencedirect.com

# **ScienceDirect**



Energy Procedia 38 (2013) 872 - 880

SiliconPV: March 25-27, 2013, Hamelin, Germany

# $Al_2O_3$ passivation on *c*-Si surfaces for low temperature solar cell applications

Desislava S. Saynova<sup>a</sup>\*, Gaby J.M. Janssen<sup>a</sup>, Antonius R. Burgers<sup>a</sup>, Agnes A. Mewe<sup>a</sup>, Elena Cianci<sup>b</sup>, Gabriele Seguini<sup>b</sup>, Michele Perego<sup>b</sup>

<sup>a</sup>ECN, Westerduinweg 3, NL-1755 LE Petten, the Netherlands <sup>b</sup>Laboratorio MDM, IMM-CNR, Via C. Olivetti 2,1-20864 Agrate Brianza (MB), Italy

#### Abstract

Functional passivation of high resistivity p-type c-Si wafer surfaces was achieved using 10 nm Al<sub>2</sub>O<sub>3</sub> layers and low temperatures for both the thermal ALD process and post-deposition anneal. Effective lifetime values higher than 1 ms were measured at excess carrier density  $\Delta n$ =10<sup>15</sup> cm<sup>-3</sup>. This result was reached in combination with temperatures of 100 °C and 200 °C for the Al<sub>2</sub>O<sub>3</sub> layer deposition and anneal, respectively. The Al<sub>2</sub>O<sub>3</sub>/c-Si interface was characterized using conductance-voltage and capacitance-voltage measurements. In particular, significantly reduced interface density of the electrically active defects  $D_{ii} \sim 2 \times 10^{10}$  eV<sup>-1</sup>cm<sup>-2</sup> was detected, which enabled excellent chemical passivation. The measured density of fixed charges at the interface,  $Q_f$ , after anneal were in the range +1 x  $10^{12}$  to -1 x  $10^{12}$  cm<sup>-2</sup> indicating that both inversion and accumulation conditions result in relevant field-effect passivation using Al<sub>2</sub>O<sub>3</sub> layers and low temperature processes. Numerical simulations on representative test structures show that the uniform  $Q_f$  effect can be understood in terms of a surface damage region (SDR) present near the interface in combination with asymmetry in the lifetime of holes and electrons in the SDR. The combination of low processing temperatures, thin layers and good passivation properties facilitate a technology for future low temperature solar cell applications.

© 2013 The Authors. Published by Elsevier Ltd.
Selection and/or peer-review under responsibility of the scientific committee of the SiliconPV 2013 conference

Keywords: Al<sub>2</sub>O<sub>3</sub> ALD layers; low temperature processing; interface passivation; numerical simulation; surface damage region

<sup>\*</sup> Corresponding author. Tel.: +31 88 515 4093; fax: +31 88 515 8214. *E-mail address:* saynova@ecn.nl.

#### 1. Introduction

Passivation schemes utilizing Al<sub>2</sub>O<sub>3</sub> ALD layers have demonstrated high performance levels on both p- and n-type c-Si surfaces of suitable doping levels for high efficiency solar cells [1, 2]. The key factor in this respect is the combination of chemical passivation enabled by low interface state defect density  $(D_{it})$ , typically 1 × 10<sup>11</sup> eV<sup>-1</sup>cm<sup>-2</sup> and field-effect passivation due to a large amount of fixed negative charges  $(Q_f)$  in the range of  $(2-6) \times 10^{12}$  cm<sup>-2</sup> located at the Al<sub>2</sub>O<sub>3</sub>/c-Si interface [3-5]. Both  $D_{it}$  and  $Q_f$  can be significantly influenced by the type of ALD, oxygen precursor, post-deposition anneal (PDA) step and film thickness [6, 7]. Thermal ALD (T-ALD) using H<sub>2</sub>O as oxygen precursor represents a particularly suitable method for low temperature applications. In particular, comparisons with plasma ALD or T-ALD with O<sub>3</sub> show that in as-deposited state typically at 200 °C the D<sub>it</sub> values using thermal ALD with H<sub>2</sub>O are lower by two orders of magnitude. As a result the passivation at that stage is better for the latter approach, even though the corresponding  $Q_f$  values are reduced by a factor of 10 [3] compared to the other two methods. This trend is also reported for PDA temperatures up to 350 °C applied for 10 min [6]. In spite of the demonstrated impressive surface passivation results with effective recombination velocities  $\leq 5$  cm/s the processing conditions are not suitable for low-temperature solar cell technologies, such as silicon heterojunction technologies. With respect to thickness, Al<sub>2</sub>O<sub>3</sub> layers deposited by thermal ALD with H<sub>2</sub>O below 10 nm result in a rapid lifetime decrease from about 1ms to only a few tens of us [6-8]. The effect is attributed to reduced chemical passivation and demonstrates the performance sensitivity towards the passivation scheme design.

This paper presents an investigation of the passivation on high resistivity p-type c-Si wafer surfaces by  $Al_2O_3$  layers with thickness of 10 nm and our focus is on low temperature processes for both thermal ALD and PDA. The characterization comprises effective lifetime measurements on symmetrical samples as well as  $D_{it}$  and  $Q_f$  evaluation using metal-oxide-semiconductor capacitors. In addition, numerical simulations on representative test structures have been carried out. The results are compared with the experimental data in order to gain a better understanding of the physical mechanism of the passivation performance and identify further improvement options for the performance of the  $Al_2O_3$  layers in low temperature solar cell applications.

# Nomenclature

 $Q_f$  fixed charges at the dielectric layer/c-Si interface

 $D_{it}$  interface state defect density

 $au_{eff}$  effective lifetime measured using photo-conductance method

 $\Delta n$  excess carrier density in the quasi-neutral region (bulk) of the silicon wafer

 $S_{eff MAX}$  maximum surface recombination velocity assuming infinite bulk lifetime,  $S_{eff} \le W/2 \tau_{eff}$ 

W wafer thickness

 $N_{acc}$  dopant concentration of p-type silicon wafer

 $S_p$ ,  $S_n$  surface recombination velocity for holes and electrons, respectively.  $S_p = \sigma_p v_{th} N_{it}$  and  $S_n = \sigma_n v_{th} N_{it}$ .

 $\tau_{s,p}$ ,  $\tau_{s,p}$  life-time of holes and electrons in the surface damage region

 $\tau_{0,p}$ ,  $\tau_{0,p}$  life-time of holes and electrons

 $\sigma_p$ ,  $\sigma_n$  capture cross section at the interface of holes and electrons, respectively

 $N_{it}$  interface defect density  $v_{th}$  thermal velocity of of the charge carriers (10<sup>7</sup> cm s<sup>-1</sup>)  $p_s$ ,  $n_s$  hole and electron concentration at the surface, respectively  $p_1$ ,  $n_1$  density of defects at midgap  $S_{eff}$  effective surface recombination velocity  $n_{b \, eff}$  effective intrinsic charge carrier density

# 2. Al<sub>2</sub>O<sub>3</sub>/c-Si interface - electrical characterization and passivation results

## 2.1. Experimental details and characterization methods

The 10-nm-thick  $Al_2O_3$  films were obtained by thermal ALD using  $Al(CH_3)_3$ , trimethyl-aluminum (TMA), and  $H_2O$  as precursors at  $T_{ALD}$  ranging between 100 °C and 200 °C. The TMA and  $H_2O$  exposure times were 0.015 s, the  $N_2$  purge time following both reactant exposures was 8 s. The growth per cycle at 100 °C was 1.0 Å/cycle, increasing to 1.1 Å/cycle at 150-200 °C. After the ALD process the films were subjected to a heat treatment in a rapid thermal anneal (RTA) furnace for 5 min in  $N_2$  atmosphere. The corresponding temperatures for this step were  $T_{PDA}$  of 200 °C or 250 °C.  $Al_2O_3$  films were deposited simultaneously on both sides of H-terminated double-side polished p-type float zone FZ<100> Si wafers (resistivity 2-4  $\Omega$  cm, thickness 280  $\pm$  25  $\mu$ m). This deposition scheme allows avoiding any difference between the two c-Si/ $Al_2O_3$  interfaces. The film thickness was measured by spectroscopic ellipsometry.

The H-terminated Si (100) surfaces were obtained by means of standard cleaning in a  $HCl:H_2O_2:H_2O=1:1:5$  solution for 10 min at 85 °C followed by a 30 s dip in diluted HF (HF:H<sub>2</sub>O=1:50) at room temperature.

Dedicated single sided deposition processes on H-terminated single-side polished p- and n- type c-Si were performed to fabricate the metal-oxide-semiconductor (MOS) capacitors for electrical characterization of the  $Al_2O_3/c$ -Si interface and evaluation of  $Q_f$  and  $D_{it}$ . The metallization of the MOS structures was obtained by evaporation technique. The front side coated with  $Al_2O_3$  was contacted via Al dots of 0.08 mm² area patterned using a shadow mask. Full Al layer deposition was applied to the opposite side. Capacitance-voltage (CV) and conductance-voltage (GV) measurements were acquired simultaneously at 100 kHz in a dark environment in a shielded probe station at room temperature. The measurements were performed from inversion to accumulation and vice versa to check the presence of hysteresis loops in the CV characteristics. Series resistance correction was applied to all measurements before the CV and GV data analysis [9].

The functionality of the surface passivation using  $Al_2O_3$  layers was monitored by effective lifetime ( $\tau_{eff}$ ) measurements on double side-coated samples as a function of the excess carrier density ( $\Delta n$ ) using a WCT-120: Silicon Wafer Lifetime Tester from Sinton Instruments in transient or quasi-steady state mode [10].

#### 2.2. Experimental results and discussion

An overview of the passivation by  $Al_2O_3$  films monitored using effective lifetime measurements and the interface properties characterized via  $Q_f$  and  $D_{it}$  evaluation is shown in Fig. 1. The reported lifetime results are averaged within each group with error bars representing the standard deviation. An experiment

group contained 2 samples. The result for each sample was averaged over 5 measurements. The plot data for  $Q_f$  and  $D_{it}$  are the average values from two different sweep directions: from inversion to accumulation and vice versa. The CV and GV measurements were of high reproducibility with essentially the same results for several capacitors and error bars comparable to the size of the symbols in the graph of Fig. 1.

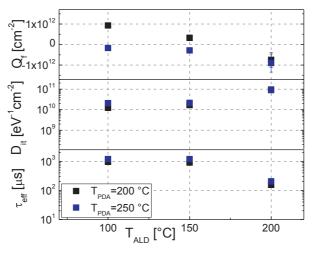


Fig. 1 Plot of the variation with the deposition temperature  $T_{ALD}$  of effective lifetime values, interface defect density  $D_{tt}$  and the fixed charge  $Q_f$ . The results shown are for two different anneal temperatures  $T_{PDA}$ 

The results show that the temperatures used for both ALD and PDA processes have an influence on the passivation performance of the Al<sub>2</sub>O<sub>3</sub> films and the characteristics of Al<sub>2</sub>O<sub>3</sub>/c-Si interface. However, the higher impact was found for the ALD conditions. Effective lifetimes  $\geq 1$  ms have been achieved even for the lowest deposition temperature of 100 °C. PDA at 250 °C shows higher and more stable results compared to 200 °C with  $(1.2\pm0.02)$  ms vs.  $(1\pm0.15)$  ms, respectively. Those values are comparable to other reports about thermal ALD Al<sub>2</sub>O<sub>3</sub> layers of 10 nm applied for passivation of (100) p-FZ wafers in the resistivity range 1.5 - 2.5 Ohm.cm [6, 7]. Nevertheless, in our case the temperatures for sample preparation were significantly reduced, both in terms of deposition (200 °C mentioned in the reports) and anneal (reference range between 350 °C and 425 °C for 10 − 15 min). In this way the overall process conditions for the Al<sub>2</sub>O<sub>3</sub> passivation layers in the current study are compatible with low-temperature solar cell technologies. The corresponding  $D_{it}$  values after anneal are  $1-2 \times 10^{10} \text{ eV}^{-1}\text{cm}^{-2}$  and significantly lower than the literature reports of  $4 - 10 \times 10^{10}$  eV<sup>-1</sup>cm<sup>-2</sup>. The chemical passivation is the most sensitive factor when the film thickness of the ALD  $Al_2O_3$  is considered [5, 7]. Higher  $T_{ALD}$  of 200 °C substantially reduced the effective lifetime to  $\sim$ 200  $\mu$ s for both anneal temperatures. Simultaneously  $D_{it}$  increased to 1  $\times$  10<sup>11</sup> eV<sup>-1</sup>cm<sup>-2</sup>. In this way the passivation was at the lowest level, although  $Q_f$ , as discussed in more details below, reached the highest negative values of  $(1 \pm 0.4) \times 10^{12}$  cm<sup>-2</sup>. This result suggests that  $D_{ii}$ values below  $1 \times 10^{11}$  eV<sup>-1</sup>cm<sup>-2</sup> are needed for the achievement of high lifetime values and the interface optimization is essential for further improvement of the low temperature passivation employing thin Al<sub>2</sub>O<sub>3</sub> layers. Concerning the field effect passivation, increasing the PDA temperature from 200 °C to 250 °C the sign of  $Q_f$  changed from positive to negative for both deposition temperatures of 100 °C and 150 °C. In case of ALD process at 200 °C the resultant  $Q_f$  was already negative and the anneal did not have a significant additional impact. The best negative  $Q_f$  of  $(1 \pm 0.4) \times 10^{12}$  cm<sup>-2</sup> was in the lower range of the typical literature data  $\sim (1-4) \times 10^{12}$  cm<sup>-2</sup>. However, higher anneal temperatures have been employed for the quoted examples. The exact mechanism of the impact on  $Q_{f}$  of the thermal treatment is not fully clarified at present and needs further investigation of the chemical and structural properties of the

Al<sub>2</sub>O<sub>3</sub>/c-Si interface. In particular, the cause for the change of the  $Q_f$  sign from positive to negative for the samples with  $T_{ALD} \le 150$  °C and  $T_{PDA} = 250$  °C can be identified by performing a depth resolved element composition quantification of the Al<sub>2</sub>O<sub>3</sub> layers. Earlier reports show that Al-rich configurations with high amount of Al<sub>i</sub> yield positive  $Q_f$ , whereas negative  $Q_f$  values correspond to O-rich structures with enhanced O<sub>i</sub> and Al vacancies [7, 11-13]. Electron energy-loss spectroscopy can also be applied to investigate Al energy-loss near-edge structures (ELNESs) and extract information about tetrahedrally and octrahedrally coordinated Al atoms in the Al<sub>2</sub>O<sub>3</sub> films [14]. The first bonding environment for Al corresponds to a negative charge while in the latter case the charge is positive [15, 16]. Another important factor influencing the  $Q_f$  in case of Al<sub>2</sub>O<sub>3</sub> layers is the interface SiO<sub>x</sub> which can be monitored using high resolution transmission electron microscopy images (HR TEM) [7, 17, 18]. In the present study we measured effective lifetime values  $\geq 1$  ms at deposition temperatures below 200 °C for both positive and negative  $Q_f$ . The best passivation in terms of lifetime values  $\sim 1.2$  ms and low spread of the results was achieved for negative  $Q_f$  and accumulation conditions.

The observation that at deposition temperatures lower than 200 °C the measured lifetime values are not strongly dependent on the sign and magnitude of the  $Q_f$  requires further investigation. The total impact of  $Q_f$  depends also on the density of the interface states and on the lifetime of the charge carriers in the wafer region close to the interface. With respect to the latter contribution it has been proposed in the literature [19, 20] that application of passivation layers may result in a so-called surface damage region (SDR) where the charge carrier lifetime is significantly lower than in the bulk of the wafer. In the next section numerical simulations are presented that include effects at the interface and in the SDR to explain why  $Q_f$  with different sign and magnitude can have a similar impact on the passivation.

#### 3. Numerical simulations

## 3.1. Approach

The modeling of the passivation properties of the  $Al_2O_3$  layers on c-Si surface is based on the well-known Shockley-Read-Hall formalism [21, 22]. The recombination at the interface is then written as:

$$R_{it} = \frac{p_s \cdot n_s - n_{i,eff}^2}{(1/S_p)(n_s + n_1) + (1/S_n)(p_s + p_1)}$$
(1)

The parameters  $S_p$  and  $S_n$  characterize the recombination at the interface defects. The charge density  $Q_f$  will have a large impact on  $n_s$  and  $p_s$ . On p-type c-Si negative  $Q_f$  will induce accumulation of holes at the interface, and positive charges will induce either depletion of the holes or even inversion conditions where  $n_s > p_s$ .

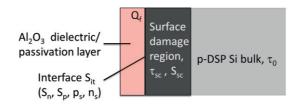


Fig. 2 Representative test structure used for the numerical simulations

The effect on the carrier concentrations in is not limited to the interface but comprises a whole space charge region (SCR), the depth of which depends on the dopant concentration of the c-Si. For doped c-Si wafers SCR is typically less than 1  $\mu$ m. As explained at the end of the previous section, in a SDR the lifetimes are reduced as compared to the bulk. Therefore, this region can contribute substantially the total recombination, and the contribution will depend on the sign and magnitude of  $Q_f$  [19, 23]. The recombination rate in the SDR is written at each point z (representing the distance to the interface) as:

$$R_{SDR}(z) = \frac{p(z) \cdot n(z) - n_{i,eff}^2}{\tau_{s,p}(n(z) + n_1) + \tau_{s,n}(p(z) + p_1)}$$
(2)

with  $\tau_{s,p}$  and  $\tau_{s,n}$  being the hole and electron lifetimes in the SDR and  $n_1$  and  $p_1$  defined as before. The passivating effect of  $Q_f$  can be characterized by  $S_{eff}$ , which is defined by:

$$S_{eff} \Delta n = R_{it}(n_s, p_s, S_n, S_p) + \int R_{SDR}(n(z), p(z) \, \tau_{s,n}, \tau_{s,p}) dz \tag{3}$$

with  $\Delta n$  being the excess carrier density directly outside the surface damage region.

Numerical simulations using the Atlas software from Silvaco [24] have been carried out to investigate the relative contributions to  $S_{eff}$  for various  $Q_f$  values. A representative test structure depicted in Fig. 2 was used. In the simulations the structure is illuminated by a monochromatic light with  $\lambda$ =500 nm.

The acceptor concentration in c-Si was fixed at  $5 \cdot 10^{15}$  cm<sup>-3</sup> corresponding to a resistivity of 2.8 Ohm cm. The SDR region had a width of 300 nm which at this resistivity is large enough to contain the space charge region. It was assumed that the lifetimes are constant over the whole SDR, i.e. in the simulations the lifetime of electrons and holes is a step function with a step at 300 nm from the interface. The effective lifetime in the SDR was adapted to the experimental results (see below). For the bulk lifetimes  $\tau_{0,p} = \tau_{0,n} = 2$  ms were assumed.

# 3.2. Simulation results and discussion

The experimental data considered for the numerical simulations correspond to  $Al_2O_3$  deposition temperatures at 100 °C and 150 °C. The selection is suitable for understanding the effect of variations in the field effect passivation, because the  $Q_f$  values change from positive to negative depending on the anneal conditions. At the same time  $D_{ii}$  remained comparable in the range of 1-2 x  $10^{10}$  eV<sup>-1</sup>cm<sup>-2</sup> providing for a high level of chemical passivation. Experimental results for the maximum surface recombination velocity  $S_{eff\ MAX}$  as function of the excess density are displayed in Figure 3a. The  $S_{eff\ MAX}$  was calculated from the measured effective lifetime values and assuming infinite bulk lifetime. The corresponding  $Q_f$  values are also included.

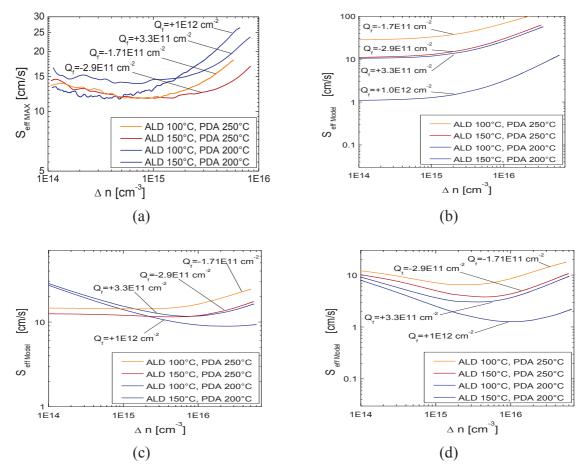


Fig. 3 (a) Experimental results for the maximum surface recombination velocity ( $S_{eff\,MAX}$ ) vs. excess carrier lifetime, (b) calculated  $S_{eff}$  vs. excess carrier lifetime without SDR, (c) numerical simulation results for  $S_{eff}$  vs. excess carrier lifetime including SDR and assuming similar lifetime values for electrons and holes, (d) model results for  $S_{eff}$  vs. excess carrier lifetime including SDR and difference by a factor of 20 in lifetime values for electrons and holes

The results of simulations without a SDR (i.e.  $\tau_{s,p} = \tau_{s,n} = 2$  ms) are shown in Fig. 3b. The limiting values in Fig. 3b at low injection are consistent with using  $S_p = S_n = 600$  cm s<sup>-1</sup>. These values were selected such that the average value of the calculated  $S_{eff}$  is in the same order of magnitude as the experimental values in Fig. 3a. Although Fig. 3b shows that  $Q_f$  with similar magnitude but opposite sign can result in the same passivation, this approach overestimates the dependence on the magnitude of  $Q_f$ .

Fig 3c shows the calculated  $S_{eff}$  when it is assumed that both the SDR and the interface contribute to the recombination. Compared to Fig. 3b the interface contribution (Eq. 1) is now reduced. i.e.  $S_p = S_n = 60$  cm s<sup>-1</sup>, and the SDR is modeled with  $\tau_{s,p} = \tau_{s,n} = 2$  µs. A good agreement could be obtained between the calculated and experimental data for the two negative charges, i.e. the cases where accumulation occurs and  $\tau_{s,n}$  determines the recombination (Fig. 3c) The positive charges induce inversion conditions and this can result in high recombination rates in the SDR especially at low injection conditions (Fig. 3c). However, the latter effect is not prominent in the experimental results (comparison between Fig. 3a and Fig. 3c). The effect of inversion conditions is mitigated by assuming  $\tau_{s,p} > \tau_{s,n}$ . Fig. 3d shows the results of

simulations using  $S_p = S_n = 60$  cm s<sup>-1</sup>,  $\tau_{s,n} = 2$  µs and  $\tau_{s,p} = 40$  µs. Comparison of Fig. 3a and Fig. 3d shows that with the assumptions made above the both the order of magnitude of  $S_{eff}$  as well as the relatively weak dependence on the sign and magnitude of  $Q_f$  can be explained.

The assumptions made here are consistent with both the results for the  $D_{ii}$  discussed above and with data from the literature. When it is assumed that the electrically active interface defects yielding the experimentally obtained  $D_{ii}$  are the main source for contribution to  $N_{ii}$ , the resultant values for the electron and hole capture cross sections are in the order of  $6 \times 10^{-16}$  cm<sup>2</sup>. This is in excellent agreement with values reported by Werner et al. for  $Al_2O_3/c$ -Si interfaces [20]. As mentioned earlier, the presence of a SDR was proposed by Steingrube et al.[19] to explain the passivation properties of  $SiN_x/c$ -Si interfaces and also used to by Werner et al. to describe the passivation at  $Al_2O_3/c$ -Si interfaces. In both papers an asymmetry  $\tau_{s,p} > \tau_{s,n}$  was assumed which was further corroborated by capture cross section measurements of Werner et al. at interface defects. It should be noted that the type of samples in our study comprising high resistivity p-type c-Si DSP wafers may facilitate higher performance sensitivity towards SDR-related effects.

An independent characterization of the SDR region is not yet available, i.e. it has not been established which type of damage has been incurred during the processing of the wafer. A tentative explanation for the presence of an SDR by Steingrube et al. is based on experimental observations that during layer deposition or during the preceding etching step hydrogen is introduced into the c-Si. High concentrations of hydrogen may then form active recombination centers. The simulation results suggest that reduction of the damage, leading to increased lifetime in the SDR while keeping the interface  $N_{it}$  at sufficiently low level, would further improve the passivation properties.

#### 4. Summary and conclusions

A high level of passivation of  $Al_2O_3/c$ -Si corresponding to  $\tau_{eff} \sim 1$  ms was achieved at low temperatures for both ALD and PDA processes (T  $\leq$  200 °C). The passivation mechanism was characterized based on  $D_{it}$  and  $Q_f$  evaluation from CV and CG measurements and numerical simulations on representative test structures. The passivation performance of  $Al_2O_3$  layers in our study was dominated by the deposition temperature rather than PDA conditions. A similar trend was found for the measured electrically active  $D_{it}$ . This is an indication of the significant impact of the chemical passivation in case of low temperature process conditions. A considerably lower influence was found for the sign and magnitude of  $Q_f$  and thereby the field effect part of the interface passivation. Numerical simulations show that the latter result cannot be understood from the recombination properties at the interface alone. The computation results strongly suggest that there is a substantial contribution from a surface damage region near the interface. This SDR is characterized by a reduced lifetime of holes and electrons, with a stronger reduction for electrons than for holes. These findings are consistent with recent literature reports. Based on the combined experimental and simulation results further improvement options were identified such as reduced layer thickness and dedicated interface optimization in order to facilitate future applications in low temperature solar cells.

## Acknowledgements

This work was supported by the EU via project NanoPV (FP7-NMP3-SL-2011-246331).

Arthur Weeber (ECN), Federico Ferrarese Lupi and Jacopo Frascaroli (MDM, IMM-CNR) are gratefully acknowledged for the fruitful discussions. Esther Cobussen-Pool (ECN) and Mario Alia (MDM, IMM-CNR) are gratefully acknowledged for the technical assistance.

#### References

- [1] J. Benick, B. Hoex, M. C. M. van de Sanden, W. M. M. Kessels, O. Schultz, S. W. Glunz, Applied Physics Letters 2008, 92, 53504
- [2] J. Schmidt, A. Merkle, R. Brendel, B. Hoex, M. C. M. van de Sanden, W. M. M. Kessels, Progress in Photovoltaics 2008, 16, 461.
- [3] G. Dingemans, N. M. Terlinden, D. Pierreux, H. B. Profijt, M. C. M. van de Sanden, W. M. M. Kessels, Electrochemical and Solid State Letters 2011, 14, H1-H4.
  - [4] B. Hoex, S. B. S. Heil, E. Langereis, M. C. M. van de Sanden, W. M. M. Kessels, Applied Physics Letters 2006, 89, 42112.
  - [5] N. M. Terlinden, G. Dingemans, M. C. M. V. de Sanden, W. M. M. Kesselsa, Applied Physics Letters 2010, 96, 12101.
- [6] G. Dingemans, R. Seguin, P. Engelhart, M. C. M. van de Sanden, W. M. M. Kessels, Physica Status Solidi-Rapid Research Letters 2010, 4, 10.
- [7] F. Werner, B. Veith, D. Zielke, L. Kuhnemund, C. Tegenkamp, M. Seibt, R. Brendel, J. Schmidt, J. Appl. Phys. 2011, 109, 13701.
- [8] J. Frascaroli, G. Seguini, E. Cianci, D. Saynova, J. van Roosmalen, M. Perego, Phys. Status Solidi A 2013, 1-5, DOI: 10.1002/pssa.201200568.
  - [9] Nicollian E.H, Brews J.R, Metal Oxide Semiconductor Physics and Technology, Wiley, New York 1982.
  - [10] A. Cuevas, R. A. Sinton, M. Kerr, D. Macdonald, H. Mackel, Solar Energy Materials and Solar Cells 2002, 71, 295.
  - [11] V. Naumann, M. Otto, R. B. Wehrspohn, C. Hagendorf, Journal of Vacuum Science & Technology A 2012, 30, D4106.
- [12] B. Shin, J. R. Weber, R. D. Long, P. K. Hurley, C. G. Van de Walle, P. C. McIntyre, Applied Physics Letters 2010, 96, 52908.
  - [13] J. R. Weber, A. Janotti, C. G. Van de Walle, J. Appl. Phys. 2011, 109, 33715.
- [14] K. Kimoto, Y. Matsui, T. Nabatame, T. Yasuda, T. Mizoguchi, I. Tanaka, A. Toriumi, Applied Physics Letters 2003, 83, 4306.
- [15] G. Agostinelli, A. Delabie, P. Vitanov, Z. Alexieva, H. F. W. Dekkers, S. De Wolf, G. Beaucarne, Solar Energy Materials and Solar Cells 2006, 90, 3438.
  - [16] R. S. Johnson, G. Lucovsky, I. Baumvol, Journal of Vacuum Science & Technology A 2001, 19, 1353.
- [17] G. Dingemans, N. M. Terlinden, M. A. Verheijen, M. C. M. van de Sanden, W. M. M. Kessels, J. Appl. Phys. 2011, 110, 93715.
  - [18] B. Hoex, J. J. H. Gielis, M. C. M. V. de Sanden, W. M. M. Kessels, J. Appl. Phys. 2008, 104, 13703.
  - [19] S. Steingrube, P. P. Altermatt, D. S. Steingrube, J. Schmidt, R. Brendel, J. Appl. Phys. 2010, 108, 014506.
  - [20] F. Werner, A. Cosceev, J. Schmidt, J. Appl. Phys. 2012, 111, 073710.
  - [21] R. N. Hall, Phys. Rev. B 1952, 87, 387.
  - [22] W. Shockley, W. T. Read, Phys. Rev. B 1952, 87, 835.
  - [23] S. Dauwe, Thesis, in , Universität Hannover 2004.
  - [24] Atlas, www.silvaco.com, 2013.